

Saved by the Bell

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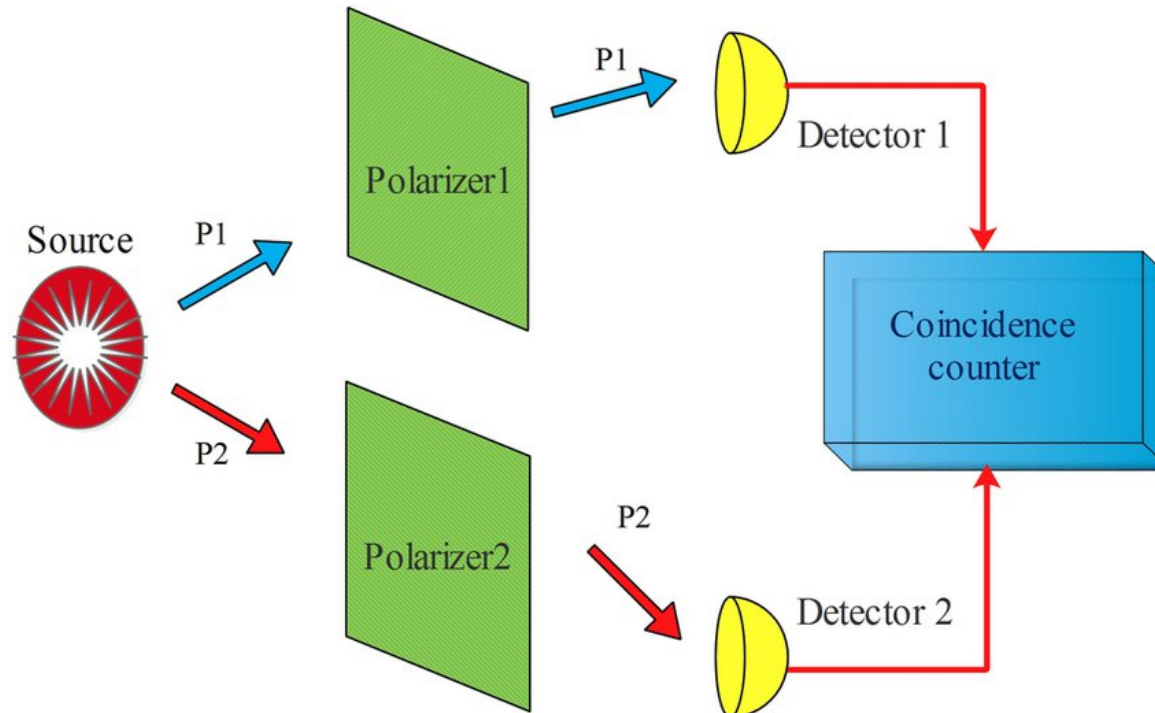
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Bell's theorem

- We start with two experimenters A and B
- Each receives a spin drawn from an entangled state
- Each measures their spin on some arbitrary axis
- They compare notes to determine the correlation between their results

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$



The result

- The resulting correlation is given as

$$C_{AB}(\varphi, \theta) = -\cos(2(\varphi - \theta))$$

- The conditional probabilities take the form

$$P(A_{\uparrow, \theta} | B_{\downarrow, \varphi}, \psi) = \cos^2(\varphi - \theta) \qquad P(A_{\uparrow, \theta} | B_{\uparrow, \varphi}, \psi) = \sin^2(\varphi - \theta)$$

- Which is taken to indicate a possible causal connection between measurements as

$$P(A_{\uparrow, \theta} | B_{\downarrow, \varphi}, \psi) \neq P(A_{\uparrow, \theta} | \psi)$$

- However, the measurements can be done at spacelike separation
- This indicates a non-local connection between outcomes

Bell's assumptions

1. Spacetime is Lorentzian
2. Causation is consistent with 1
3. Probability completeness

$$P(Y) = \sum_X P(X)P(Y|X)$$

4. Common cause screening

$$P(A, B|X) > P(A|X)P(B|X)$$



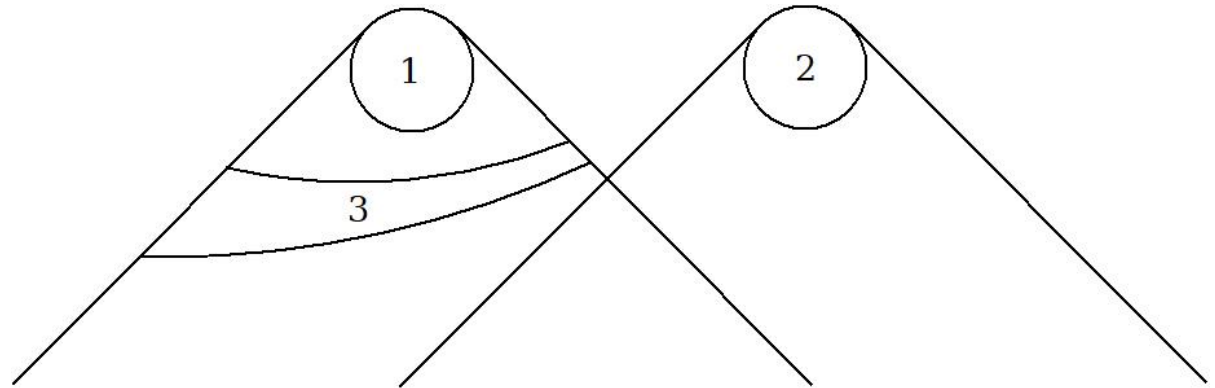
A causes B
B causes A
Something in X causes both

5. Probability factorisation

$$P(A|B, X) = P(A|X)$$



If 4 is true, X can contain a common cause



Which assumption fails?

- In quantum mechanics it is probability factorisation

$$P(A_{\uparrow,\theta}|B_{\downarrow,\varphi},\psi) = \cos^2(\varphi - \theta) \quad P(A_{\uparrow,\theta}|B_{\downarrow,\varphi},\psi) \neq P(A_{\uparrow,\theta}|\psi)$$

- However, this can easily fail in classical cases

A machine deals out 2 coins
showing opposite faces.

$$P(A = 1|B = 0) = 1 = P(B = 1|A = 0)$$

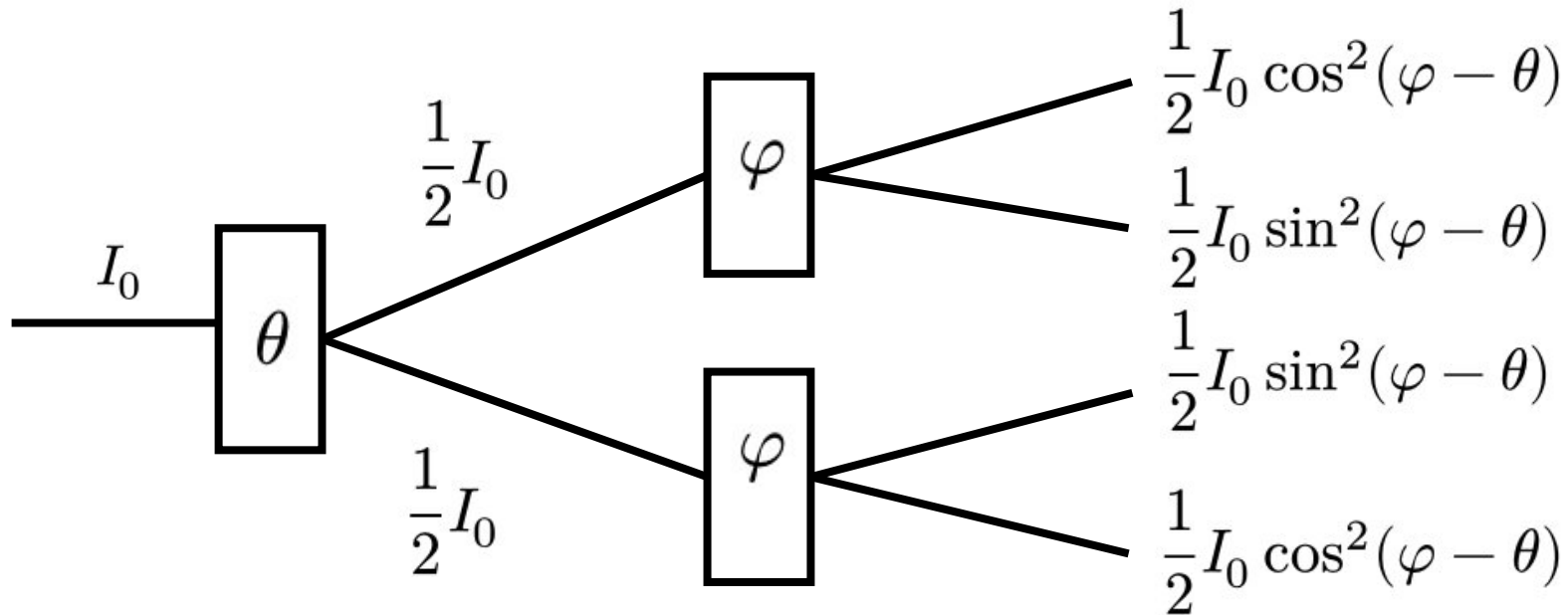
$$P(B = 0) = \frac{1}{2} = P(A = 0)$$

Implies causal link
between A and B,
despite existing
common cause

$$P(A = 1, B = 0) = \frac{1}{2} > P(A = 1)P(B = 0)$$

A classical polarisation test

- We start with an unpolarised state
- Pass through two layers of polarising beam splitter



- The correlation function matches Bell up to a sign
- Is this a coincidence?

$$C(\varphi, \theta) = \cos(2(\varphi - \theta))$$

Coincidence or not?

- Our Bell state particles are highly correlated
- If we measure both on the same axis we get perfect matching/anti-matching

$$P(A_{\uparrow,\theta}|B_{\downarrow,\theta}, \psi) = 1$$

- This means our conditional probabilities obey

$$P(A_{\uparrow,\theta}|B_{\downarrow,\varphi}, \psi) = P(A_{\uparrow,\theta}|A_{\uparrow,\varphi}, \psi)$$

- What is this probability?
- It is that of successive measurement results on spin A alone
- Interestingly,

$$\text{Tr}(\hat{\rho}\hat{\sigma}_{\theta}\hat{\sigma}_{\varphi}) = \cos(2(\varphi - \theta)) = \text{Tr}(\hat{\rho}\hat{\sigma}_{\varphi}\hat{\sigma}_{\theta})$$

- **Our successive and entangled experiments are two versions of the same measurement**

What does it imply?

- There is nothing non-local/entangled about our successive experiment
- The results are just from the physics of filtering spins
- **Could this also be true for the entangled mirror case?**
- It is odd that the correlation is the same with different explanations
 - Entangled case: non-locality
 - Successive case: local physics
- Could we use the matching correlation to explain the physics?

Directional asymmetry

- The entangled case can't explain the successive one
- However, the same local physics is present in both cases
- Says that measuring A tells us the B value on the same axis
- This is a counterfactual
- The two experiments are counterfactual mirrors
- Quantum mechanics is “counterfactually definite”

$$P(A_{\uparrow,\theta}|B_{\downarrow,\varphi},\psi) = P(A_{\uparrow,\theta}|A_{\uparrow,\varphi},\psi)$$

Measuring B tells me the unmeasured response of A!

Counterfactuals and non-locality

- The entanglement allows us to counterfactually measure B via A
 - This is why the correlation is the same
 - Entanglement doesn't add extra physics
 - It allows us to test the same physics differently
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- This explains why our “non-locality” violates the spirit of relativity but not the letter
 - It isn't actually non-local



If there's no non-locality, why don't local hidden variable models work?

The LHV problem

- Let's take a look at how these models are often framed

$$S_{\theta} = \text{sgn}(\vec{\theta} \cdot \vec{\lambda})$$

- This yields discrete results.
- It is still classical
- It is assumed there is a vector description of spin

Bell, 1965

The alternative?

- Spins have the same “projection” on every axis
- This cannot be encompassed by a vector
- Quantisation of the spin is what prevents this
- It's possible that what hidden variable models are missing is quantisation
- Of course, it's not clear how we could incorporate this mathematically!

A Fine solution

- Fine's theorem says that the non-existence of

$$P(s_1, s_2, s_3)$$

- Is both necessary and sufficient for Bell violation
- This makes no sense in the context of Bell experiments!
- These deal in two-spin correlations, not three!
- In a successive scenario it just means an extra layer of filters
- Open question: does the third layer requirement reveal any physics?

Fine, 1982

Halliwel, 2014 arXiv: 1403.7136

Conclusions



- There is no coincidence in the successive-entangled matching
 - They are counterfactual equivalents
 - This requires a realist belief in premeasurement values
 - Bell experiments are local in this view
 - Local hidden variables still fail
 - Just due to not taking quantisation seriously?
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- Bell says that no local, **classical**, hidden variable model can reproduce QM